

Data-Driven Review of Thermoelectric Materials: Performance and Resource Considerations

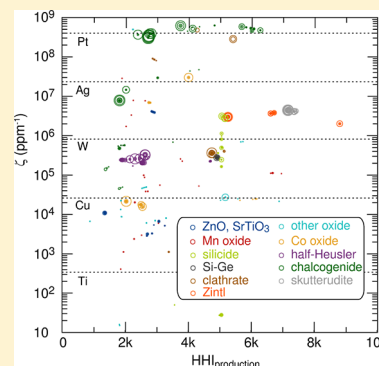
Michael W. Gaultois,^{*,†,‡} Taylor D. Sparks,^{*,‡} Christopher K. H. Borg,[‡] Ram Seshadri,^{*,§,†,‡} William D. Bonificio,^{||} and David R. Clarke^{||}

[†]Department of Chemistry and Biochemistry, [‡]Materials Research Laboratory, and [§]Materials Department, University of California, Santa Barbara, California 93106, United States

^{||}School of Engineering and Applied Sciences, Harvard University, 29 Oxford Street, Cambridge, Massachusetts 02138, United States

ABSTRACT: In this review, we describe the creation of a large database of thermoelectric materials prepared by abstracting information from over 100 publications. The database has over 18 000 data points from multiple classes of compounds, whose relevant properties have been measured at several temperatures. Appropriate visualization of the data immediately allows certain insights to be gained with regard to the property space of plausible thermoelectric materials. Of particular note is that any candidate material needs to display an electrical resistivity value that is close to 1 m Ω cm at 300 K, that is, samples should be significantly more conductive than the Mott minimum metallic conductivity. The Herfindahl–Hirschman index, a commonly accepted measure of market concentration, has been calculated from geological data (known elemental reserves) and geopolitical data (elemental production) for much of the periodic table. The visualization strategy employed here allows rapid sorting of thermoelectric compositions with respect to important issues of elemental scarcity and supply risk.

KEYWORDS: thermoelectrics, datamining, Herfindahl–Hirschman Index, elemental abundance



INTRODUCTION

The nature of thermoelectric phenomena and materials—competing and contraindicated properties, the complexity and variety of the material systems involved—make it somewhat difficult to develop rational strategies that can lead to significant improvements in performance. Notwithstanding these difficulties, creative approaches have yielded highly promising materials.^{1–8} The guiding principle behind the design of thermoelectric materials, and indeed, any functional material, is to completely understand the causal physics and use such knowledge to rationally optimize material properties. However, even without knowledge of causality, progress across numerous fields is enabled by correlation, without a priori understanding of the drivers. Particularly when large data sets are available, robust correlation can be found between seemingly disparate observables. This ability to extract meaningful information from large pools of data has been somewhat under-utilized in the search for new materials.

Thermoelectrics and their historical development have been surveyed by Snyder and Toberer,⁹ Nolas et al.,¹⁰ Goldsmid,¹¹ and the two CRC handbooks on thermoelectrics by Rowe.^{12,13} The volume and complexity of research on thermoelectric materials makes the field fertile for a data-driven review—sometimes referred to as data mining, or materials informatics. Informatics-based approaches have been successfully used for estimating some physical properties¹⁴ and the relative stability of selected material systems,¹⁵ and high-throughput methods are becoming increasingly helpful in materials design.¹⁶ Given the volume of information that has been published, knowing

the right information to abstract is the first step to such an approach. Furthermore, developing an appropriate visualization strategy to explore the space of thermoelectric materials is crucial. Accordingly, we have reduced the problem of reviewing thermoelectric performance to several key properties at four temperatures of interest, and have created an interactive framework to visualize the large amounts of information. Trends in materials properties emerge from such visualization, and lead to guiding principles for the development of high-performance thermoelectric materials. Perhaps more importantly, guidelines are suggested for where *not* to look in the parameter space of candidate materials.

Since the efficiency of a thermoelectric device is related to the thermoelectric figure of merit, $zT = S^2T/(\rho\kappa)$, with S the Seebeck coefficient, ρ the electrical resistivity, and κ the thermal conductivity, these different measured parameters were the obvious choice for extraction from publications into a database. A particularly useful format for displaying the large amount of property-based data is based on a modification of the well-known Jonker plot,¹⁷ in which the Seebeck coefficient is usually graphed as a function of the electrical conductivity. The use of the Jonker plot has traditionally been limited to optimizing the carrier concentration in a single thermoelectric material. In contrast, we display *all* materials in a single plot, with the different material classes grouped by marker color. In the

Received: March 18, 2013

Revised: May 2, 2013

Published: May 6, 2013



modified version employed by us, we prefer to use the electrical resistivity (on a logarithmic scale) ρ as abscissa, and the Seebeck coefficient S as ordinate. The radius of the circular marker represents the performance, which is commonly either zT or the power factor, S^2/ρ .

There is more to a material than just performance. Given the proposed widespread application of thermoelectrics and the potential for high-volume use of materials if certain performance targets are reached, they must be composed of elements that are accessible and not in danger of a supply risk. Crustal abundance of elements, their global production, reserves, and use, are some of the factors that determine supply risk. The criticality of elements in the context of metals that are crucial to energy conversion has been described by Graedel.¹⁸ Homm and Klar¹⁹ have specifically raised these issues for thermoelectric materials. Following the approach of Graedel,¹⁸ we have used our database to calculate several criticality indices for the thermoelectric materials featured here. Resources that are produced almost entirely in a particular region can provide a single entity leverage in determining supplies and prices. This geopolitical influence over materials supply and price can be measured by market concentration, often quantified through the Herfindahl–Hirschman Index (HHI).^{20,21} The HHI is a financial tool commonly used to measure the monopoly of entities over a commodity or product, and has been previously used as a measure of geopolitical influence on elemental production of a select few elements.^{18,22,23} Here we calculate the HHI, based on available 2011 data, of almost all of the first 83 elements in the periodic table. Only H, the noble gases, Tc, and Pm are excluded. The HHI is calculated both for elemental production (HHI_p), reflecting the specific geopolitics of the element, as well as for elemental reserves (HHI_R), based on known deposits that could be processed. In conjunction with composition data entered into the database, we employ HHI indices and elemental scarcity values to determine practical issues that influence the likelihood a particular thermoelectric material will find widespread use.

METHODS

Nature and Source of Data. We have abstracted data from representative publications on a wide variety of thermoelectric compounds so that they can be accessed and compared easily with other compounds. The data correspond to four different temperatures in the different regimes of interest for high-temperature thermoelectric devices: 300 K, 400 K, 700 K, 1000 K. There are currently over 1100 database entries (rows), each with 17 associated components: temperature, electrical resistivity, Seebeck coefficient, thermal conductivity, power factor, κzT , zT , chemical composition, material family, preparatory route, material form (whether single crystal or polycrystalline), author, year of publication, DOI link, and comments. Additionally, we generate new metadata from the chemical composition of materials, such as HHI_p, HHI_R, scarcity, and average atomic weight, \bar{M} . To facilitate comparison of materials, we limit our choice of systems in the database to be bulk, polycrystalline samples. In many instances, we employ the power factor or κzT rather than the actual figure of merit zT for the visualization, because zT requires κ , the thermal conductivity, to be measured. Thermal conductivity is not reported for many materials. Furthermore, of all the thermoelectric parameters it is the least reliable and/or reproducible because it is so sensitive to processing conditions.

Data from published work was extracted manually from digital publications using free software such as PlotDigitizer²⁴ and DataThief.²⁵ In general, data was extracted from plots of a physical property vs temperature. If data was not explicitly reported at a temperature of interest, values were interpolated, or extrapolated when appropriate. If

property traces (curves) were found to be rapidly or unpredictably changing in the region of extrapolation, the point was omitted, or data was taken from the nearest reported temperature. In these cases, the temperature of an extrapolated data point is mentioned in the metadata comment. Because data was entered by hand, the power factor and figure of merit (zT) were calculated from the extracted data and checked against the reported values to ensure the data was self-consistent.

HHI values based on production and reserves for each element were calculated from 2011 USGS commodity statistics following the approach used by others.^{18,22,23} When 2011 data was unavailable, data from 2010 or 2009 was used. This has generated a set of HHI values presented in Table 1 for much of the periodic table, a significant expansion from previous studies, which focused on eight parent metals.^{18,23} For elements where reserves are seen as adequate or extremely large (e.g., C, O, F, Na, Al, Ca, S, etc.) quantitative reserves are not available, but their use is unlikely to reduce their availability to critical levels. For elements such as sodium, where multiple

Table 1. Herfindahl–Hirschman Index (HHI) Calculated for Much of the Periodic Table, Using Recent USGS Data, Typically 2011 for Most Materials, But No Older than 2009^a

element	HHI _p	HHI _R	element	HHI _p	HHI _R
He	3200	3900	Ru	3200*	8000*
Li	2900	4200	Rh	3200*	8000*
Be	8000	4000*	Pd	3200	8000*
B	2900	2000	Ag	1200	1400
C	500*	500*	Cd	1700	1300
N	1300	500*	In	3300	2000*
O	500*	500*	Sn	2600	1600
F	1500*	1500*	Sb	7900	3400
Na	1100	500*	Te	2900	4900
Mg	5300	500*	I	4900	4800
Al	1600	1000*	Cs	6000*	6000*
Si	4700	1000*	Ba	3000	2300
P	2000	5100	La	9500	3100
S	700	1000*	Ce	9500	3100
Cl	1500*	1500*	Pr	9500	3100
K	1700	7200	Nd	9500	3100
Ca	3900	1500*	<i>Pm</i>	9500	3100
Sc	5500*	4500*	Sm	9500	3100
Ti	1100	1600	Eu	9500	3100
V	3300	3400	Gd	9500	3100
Cr	3100	4100	Tb	9500	3100
Mn	1600	1800	Dy	9500	3100
Fe	2400	1400	Ho	9500	3100
Co	3100	2700	Er	9500	3100
Ni	1000	1500	Tm	9500	3100
Cu	1600	1500	Yb	9500	3100
Zn	1600	1900	Lu	9500	3100
Ga	5500*	1900*	Hf	3400*	2600*
Ge	5300	1900*	Ta	2300	4800
As	3300	4000*	W	7000	4300
Se	2200	1900	Re	3300	3300
Br	3300	6900	Os	5500*	9100*
Rb	6000*	6000*	Ir	5500*	9100*
Sr	4200	3000*	Pt	5500	9100*
Y	9800	2600	Au	1100	1000
Zr	3400	2600	Hg	5500	3100
Nb	8500	8800	Tl	6500*	6500*
Mo	2400	5300	Pb	2700	1800
			Bi	5300	6000

^aAsterisks against numbers indicate some uncertainty (see text).

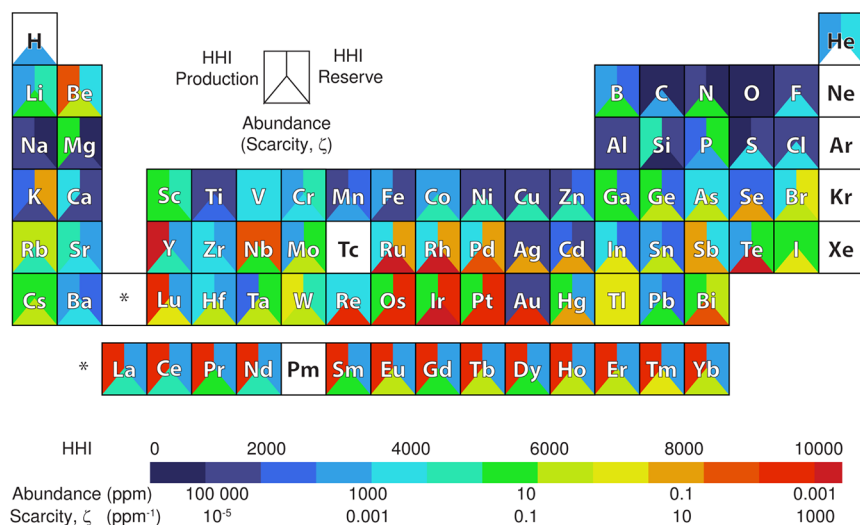


Figure 1. Periodic table indicating elemental scarcity and the HHI (production, reserves) indices for most elements. Values were calculated using USGS statistics from 2011, or if unavailable, from statistics no earlier than 2009.

commodities are reported separately (carbonates, sulphates, chloride ...), the reports were combined and the HHI values were generated from the aggregate. Another important consideration is that the production and/or reserve values for a country may be unknown or are withheld. In these instances, where the number of producers and/or the amount of reserves are low, these omissions introduce considerable uncertainty. In these cases, the HHI values were estimated based on the general information provided in the USGS report. Estimates of these elements are denoted by an asterisk in Table 1.

HHI values were calculated using the following expression, $HHI = \sum_i^N s_i^2$, where N is the total number of countries involved, and s_i is the percent market share of country i in the world production or reserves of a given element. The U.S. Department of Justice and the Federal Trade Commission have designated markets as unconcentrated when $HHI < 1500$, moderately concentrated when the HHI lies between 1500 and 2500, and highly concentrated when $HHI > 2500$.²⁶ If a single country controlled the entire market, $HHI = 100^2$. Elemental HHI values were then used to calculate weighted HHI production and reserve values based on the *weight* fraction of each element in the chemical formula. Atomic weights were taken from the CRC Handbook.²⁷ The crustal abundance of elements were obtained from the CRC Handbook²⁷ and used to generate elemental scarcity values, ζ (crustal abundance in inverse ppm). These scarcity values were used to calculate the effective scarcity of materials based on the weight fraction of elements in the chemical formula, $\zeta = \sum_i^N (\zeta_i \times (m_i/m_{tot}))$ where ζ_i and m_i are the scarcity and weight of an element i in a material that has N different elements. The scarcity and Herfindahl-Hirschman Index values (based on production and reserves) for much of the periodic table is shown schematically in Figure 1.

Mechanics of Visualization. At the initial stages of this work, it became evident that gathering the large amount of data would be futile in the absence of an appropriate framework that permitted the gathered data to be visualized in a flexible manner. In this section we describe a Web site we have developed and hosted at <http://www.mrl.ucsb.edu:8080/datamine/thermoelectrics.jsp> as an example of a plausible framework for organizing and visualizing the results of such data mining. The flowchart describing the data mining, database formation, and visualization process is summarized in Figure 2. A screenshot of the Web site is shown in Figure 3. Because several physical parameters were collected and tabulated, visualizing any number of combinations along the abscissa and ordinate is possible using the Web site. While not all combinations will yield insightful relationships, the Web site allows using as axes choices the electrical resistivity, Seebeck coefficient, thermal conductivity, average atomic

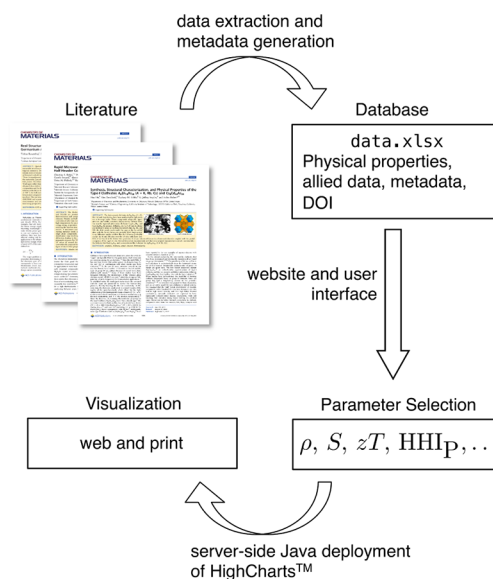


Figure 2. Flowchart for extraction, organization, and visualization of data. Note that the first step of data extraction is manual.

weight, elemental scarcity, and HHI based on production or reserves of a material. To increase the information density that can be visualized, a third dimension is also plotted: the size of the data point (the radius of a circle) is proportional to material performance. The marker size can be either power factor (S^2/ρ), κzT , or zT . κzT allows a rough comparison with zT , but does not require the thermal conductivity to have been measured or reported. Many materials have a thermal conductivity between $1 \text{ W m}^{-1}\text{K}^{-1}$ to $10 \text{ W m}^{-1}\text{K}^{-1}$, and in the case that $\kappa = 1 \text{ W m}^{-1}\text{K}^{-1}$, the numerical value of κzT is the same as the numerical value of zT .

The Web site allows interactive exploration of all the data. Hovering over a data point reveals a tooltip with pertinent information: the values for the abscissa and ordinate, the chemical composition, the sample form (e.g., polycrystalline, single crystal, nanoparticles), the preparatory route (e.g., ceramic method, arc melting), the author and year, and either the power factor, κzT , or zT value. Additionally, clicking on a data point leads the web-browser directly to the publication via the document object identifier (DOI). To enable sorting of the large number of data sets, the user may choose to sort by material family or temperature regime. Although the material families

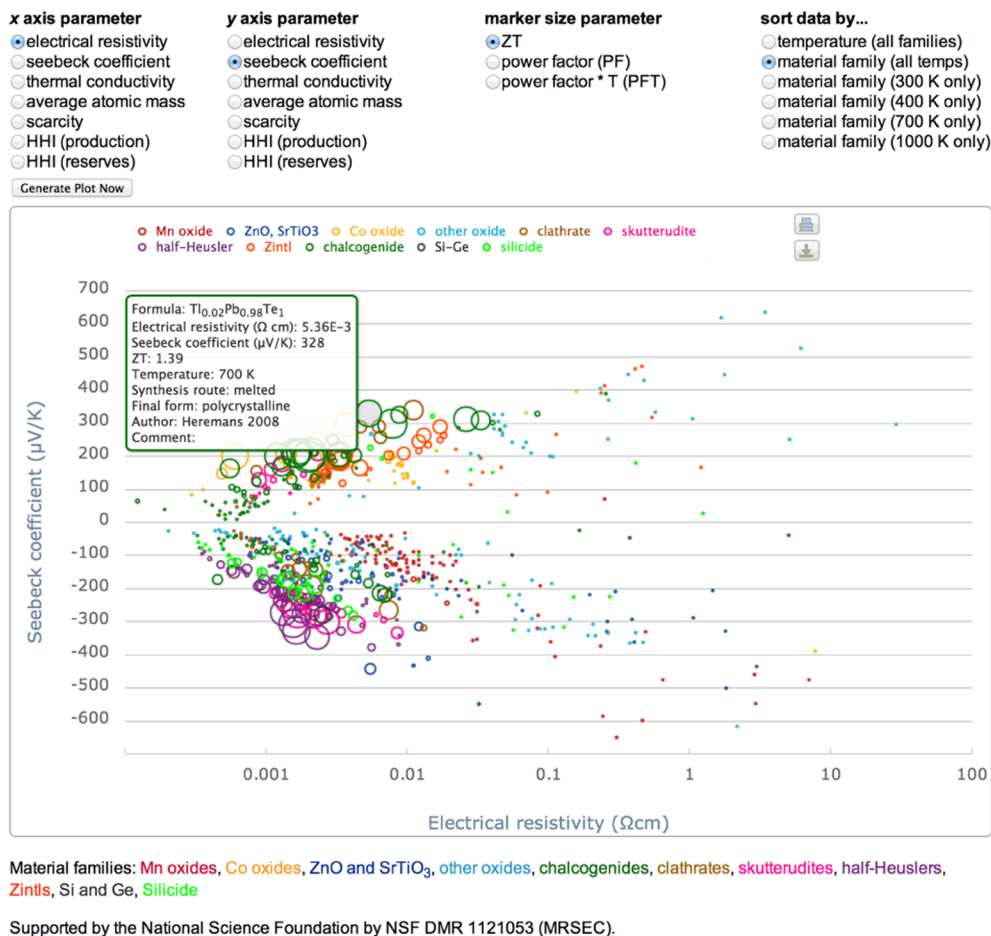


Figure 3. Screenshot of the web-based visualization tool, that permits the simultaneous visualization of four parameters: abscissa, ordinate, marker size, and color: Several variables can be chosen as abscissa and ordinate, and measures of thermoelectric performance can be represented by the radius of the data points. To simplify navigation, families of related materials can be displayed or hidden by clicking their legend marker. Further, hovering over a data point intuitively reveals a tooltip with pertinent information: the names and values for the abscissa and ordinate, the chemical composition, the sample form (e.g., polycrystalline, single crystal, nanoparticles), the preparatory route (e.g., ceramic method, arc melting), the author and year, and either the power factor, κzT , or zT value. Clicking on a data point leads the web-browser directly to the source publication via the document object identifier (DOI).

Table 2. Flexibility of the Visualization Framework Allows Users to Examine the Relation between Any Number of Different Parameters^a

abscissa	ordinate	size (radius)	use	finding
ρ	S	$S^2/\rho, \kappa zT, zT$	general trends, insight	Compound must be metallic for high performance
\bar{M}	κ	$\kappa zT, zT$	effect of \bar{w}	κ decreases with increasing \bar{M}
$\rho_{300\text{ K}}$	$S_{300\text{ K}}$	$\kappa zT, zT$ at different temperatures	estimate high- T performance	Rapid screening of high- T materials
ρ	κ	$\kappa zT, zT$	effect of κ_c	Slope of κ vs ρ changes linearly with T of data set
HHI_p	HHI_r	$\kappa zT, zT$	material choice and criticality	Intensive use of rare-earths and/or Sb may strain markets
ζ	HHI_p	$\kappa zT, zT$	material choice and criticality	Some elements are abundant, yet will have volatile prices. Many state-of-the-art high-performance materials rank poorly in criticality indices

^aSeveral examples are tabulated here, along with what can potentially be learned.

are shown with different marker color the ability to hide or show a given family can make a direct comparison more clear. This is accomplished by clicking the name of the data set in the legend. Finally, option is available to visualize the results of the thermoelectric database employed here, or for users to upload their own data, following an Excel template file available on the Web site. In this way a user can look for trends in their data or use the Web site's Java code to generate additional data not related to performance such as scarcity, average atomic weight, or HHI.

RESULTS AND DISCUSSION

Examples of Useful Visualization Schemes. In Table 2, we list those combinations of parameters that we have found particularly useful to plot. The various combinations and the findings are described more thoroughly in the following sections.

Rapid Screening of Materials. We begin by visualizing the general trend in Seebeck coefficient as a function of the

electrical resistivity (Figure 4): These two properties jointly contribute to zT . Higher information density is achieved by

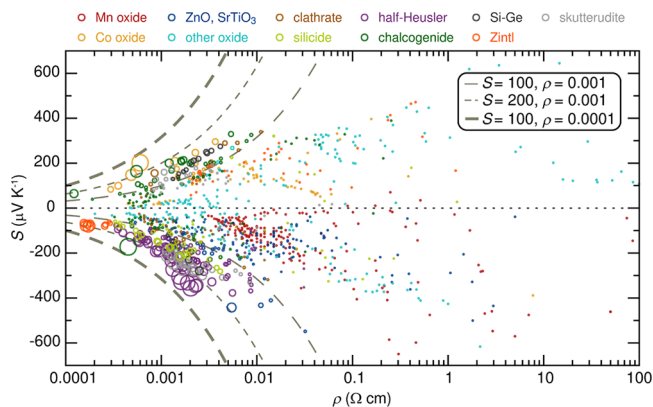


Figure 4. Seebeck coefficient of a wide variety of materials grouped by material family is plotted against electrical resistivity. The marker size (radius) here is proportional to the power factor. Materials with high performance exhibit a high Seebeck coefficient at a given electrical resistivity, and lie outside the conical envelope defined by most thermoelectric materials. This envelope can be described by lines of constant power factor (S^2/ρ), which are shown to give lower and upper bounds of performance. With few exceptions, all investigated materials with reasonable thermoelectric performance are well described as metals ($\rho_{300\text{ K}} \ll 0.01\ \Omega\ \text{cm}$).

encoding the power factor as the marker size. This style of plot was first introduced in a previous report.²⁸ Here we enhance it by assigning a color to a family of related materials. Although the axes are similar to those for a Jonker plot (where Seebeck coefficient is plotted vs electrical conductivity) the application is quite distinct. Jonker plots are usually employed to examine the effect of changing the carrier concentration in a single material.^{17,29}

Looking at the data as a whole, the general envelope of materials is cone-like, widening at higher electrical resistivities (Figure 4). The best materials clearly define the left-most edges of the envelope. All the high-performance materials have interesting physics that make them special, even when looking only at these two properties. Specifically, materials with metallic behavior generally have low Seebeck coefficients, but high-performance thermoelectric materials violate this principle. For example, band asymmetry or high band degeneracy near the Fermi level in BiTe- and PbTe-based systems leads to an unusually high Seebeck coefficient, despite their metallic behavior.^{5,30} Likewise, even though Na_xCoO_2 is metallic, with polycrystalline samples having $\rho_{300\text{ K}} \approx 0.002\ \Omega\ \text{cm}$, the correlated behavior of electrons and spin contribution to thermopower—or, arguably, the unique band structure—lead to a remarkably high Seebeck coefficient.^{31,32}

Looking at the ensemble of data (Figure 4) also provides insight about particular families of materials. For example, we find that optimizing nonmetallic moderate-performance materials is unlikely to lead to dramatic gains in performance, as the electrical resistivity is too high. For example, CaMnO_3 -based systems remain a topic of intense study, but when viewed in the context of all materials, their placement in the map suggests it is unlikely that high performance will ever be reached.

It was previously shown that room-temperature thermoelectric properties (Seebeck coefficient and electrical resistivity)

of several high-performance thermoelectric materials clustered in one area of the thermoelectric map (Figure 4) when compared with low-performance materials.²⁸ Here we examine the relationship between the properties at room-temperature and the properties at the temperature of highest zT for a small set of mid- and high-temperature thermoelectric materials (Figure 5). Interestingly, the properties of several different

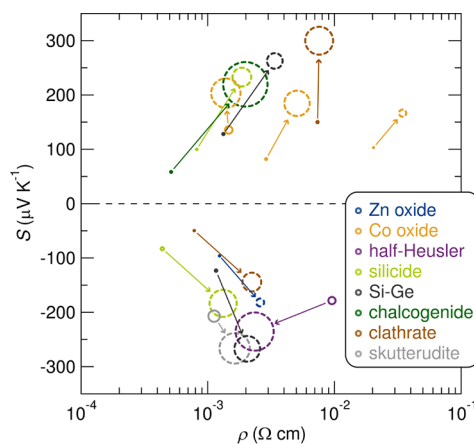


Figure 5. Several high-performance thermoelectric materials were chosen to examine how their room temperature properties relate to their properties at the temperature of highest zT . Arrows show the change from room temperature (solid circles) to the temperature of highest zT (dashed circles). The properties of several different material families appear to cluster at their temperature of highest zT , with the electrical resistivity between $0.001\ \Omega\ \text{cm}$ and $0.01\ \Omega\ \text{cm}$ and the absolute value of the Seebeck coefficient between $150\ \mu\text{V}\ \text{K}^{-1}$ and $300\ \mu\text{V}\ \text{K}^{-1}$. Additionally, nearly all high-performance materials have a room-temperature electrical resistivity below the Mott maximum metallic resistivity ($\rho_{300\text{ K}} < 0.01\ \Omega\ \text{cm}$). The materials shown are $\text{Zn}_{0.98}\text{Al}_{0.02}\text{O}$,³³ Na_xCoO_2 ,³⁴ $\text{Ca}_2\text{Co}_2\text{O}_5$,³⁵ $(\text{Zr}_{0.5}\text{Hf}_{0.5})_{0.5}\text{Ti}_{0.5}\text{NiSn}$,³⁶ $\text{Mg}_2\text{Si}_{0.999}\text{Bi}_{0.001}$ (*p*-type),³⁷ $(\text{Mg}_2\text{Si})_{0.97}\text{Bi}_{0.03}$ (*n*-type),³⁸ $\text{Si}_{0.8}\text{Ge}_{0.2}$ (*p*-type),³⁹ $\text{Si}_{0.8}\text{Ge}_{0.2}$ (*n*-type),³⁹ $\text{PbTe}_{0.75}\text{Se}_{0.25}$,³ CsBi_4Te_6 ,⁴⁰ $\text{Ba}_8\text{Ga}_{18}\text{Ge}_{28}$ (*p*-type),⁴¹ $\text{Ba}_8\text{Ga}_{16}\text{Ge}_{30}$ (*n*-type),⁴² and $\text{In}_{0.25}\text{Co}_4\text{Sb}_{12}$.⁴³

material families appear to cluster at their temperature of highest zT , with the electrical resistivity between $0.001\ \Omega\ \text{cm}$ and $0.01\ \Omega\ \text{cm}$ and the absolute value of the Seebeck coefficient between $150\ \mu\text{V}\ \text{K}^{-1}$ and $300\ \mu\text{V}\ \text{K}^{-1}$. Although the small data set prevents any strong conclusions, examining more materials may provide some predictive ability of high-temperature properties, using those measured at room temperature. Even some hints with regard to high-temperature properties would be powerful: measurement of high temperature properties is time-consuming and requires specialized instruments. Furthermore, there is much more room-temperature data available in the literature, and a single room-temperature measurement could facilitate combinatorial testing of a large phase space.⁴⁴

The thermoelectric survey shown in Figure 4 and the high-temperature trends observed in Figure 5 suggest that all thermoelectric materials with any appreciable performance are all metallic, with an electrical resistivity well below the Mott maximum metallic resistivity at room temperature, that is, $\rho_{300\text{ K}} < 0.01\ \Omega\ \text{cm}$. This provides a valuable thermoelectric screening criterion and guiding direction for future studies. Unlike intermetallic compounds, which are generally metallic or semimetallic, transition metal oxides span the gamut of electrical resistivities, ranging from insulating (e.g., TiO_2) to metallic (e.g., ReO_3). A common strategy to seeking effective thermoelectric oxides is to examine the metal/nonmetal border,

that is, proximal to the Mott minimum metallic conductivity, where Seebeck coefficients are often substantial. However, the utility of this approach appears somewhat questionable, given that the good thermoelectrics appear well on the metallic side.

Thermal Conductivity. Over the past several decades, the majority of improvements in thermoelectric materials have resulted from decreasing the lattice contribution to thermal conductivity. A number of materials selection guidelines have been identified to reduce the thermal conductivity in a material.^{9,45,46} One such strategy is to lower the vibrational frequency, and thus the thermal conductivity, by using materials with a large average atomic weight. This has frequently been touted as a reason for the higher thermal conductivity in oxides relative to compounds with heavier anions such as Te, Se, and Sb. Examination of the relationship between thermal conductivity, κ , and average atomic weight, \bar{M} , confirms the general reduction of thermal conductivity in heavier compounds (Figure 6). For example, compounds with $\bar{M} \approx 25$ g/

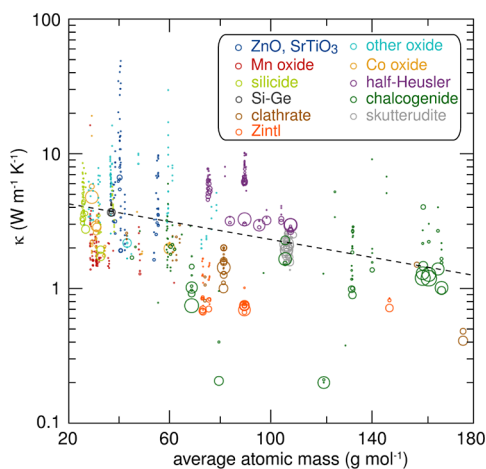


Figure 6. Thermal conductivity plotted against average atomic weight for a variety of thermoelectric materials grouped by material family, with marker indicating zT . The dashed line represents a best fit regression. Notwithstanding the combination of pure phase materials and heavily nanostructured materials, the thermal conductivity generally decreases with increasing average atomic weight. However, at fixed average atomic weight, there is considerable tunability.

mol have an average $\kappa \approx 4 \text{ W m}^{-1} \text{ K}^{-1}$ whereas those with $\bar{M} \approx 105$ g/mol have an average $\kappa \approx 2 \text{ W m}^{-1} \text{ K}^{-1}$. In the absence of point defects or other scattering mechanisms, the reduction in thermal conductivity should scale as $\bar{M}^{-1/2}$.⁴⁷ However, many of the materials contained in this plot rely on additional techniques to reduce thermal conductivity (e.g., increased phonon scattering from defects, alloying, grain boundaries, interfaces, nanobulk compounds, complex crystal structures, etc.). On the other hand, the lower vibrational frequency achieved in compounds with large average atomic weight has only a small impact on the electrical resistivity.

Finally, we examine the relationship between total thermal conductivity and electrical resistivity (Figure 7); the thermal conductivity decreases with increasing electrical resistivity. This behavior is expected for the electronic contribution from the Wiedemann–Franz law, $\kappa_e/\sigma = \pi^2 k_B^2/3e^2 T = L_0 T$ where L_0 is the Lorenz number ($2.44 \times 10^{-8} \text{ W } \Omega \text{ K}^{-2}$) and σ is the electrical conductivity. There is another consideration that could be important. For a given composition, the same structural features that give rise to low resistivity, for example

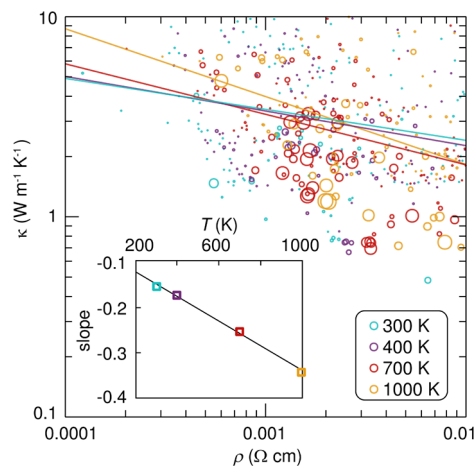


Figure 7. Thermal conductivity plotted against electrical resistivity for a variety of thermoelectric materials grouped by data set temperature with marker size indicating zT . Lines are best fit linear regressions of the data plotted on log–log plot. The slopes are then plotted as a function of data set temperature in the inset. Despite the large scatter in the data that has been fitted, the general trend that emerges is consistent with the T^{-1} dependence of Umklapp scattering of phonons.

structures that are highly connected in three-dimensions, are the same features that often lead to high lattice thermal conductivity.

Examining the broad correlation across all material families in the metallic region ($\rho < 0.01 \text{ } \Omega \text{ cm}$) reveals the slope of the log–log plot κ vs ρ decreases linearly with increasing temperature (inset, Figure 7). The power-law behavior revealed from the fit is consistent with Umklapp scattering, explained below.

Examining correlations across many material families with disparate properties will only provide general trends rather than precise values for a specific material family (Table 3).

Table 3. General Class of Material Family and Publications from Which Data Was Extracted^a

material family	references
Mn oxide	49–60
ZnO, SrTiO ₃	61–72
Co oxide	73–77
other oxide	28,63,78–96
Si–Ge	39,97
clathrate	42,42,98–103
half-Heusler	36,104–107
skutterudite	1,7,43,108
chalcogenide	5,30,40,109–122
silicide	37,38,123–126
Zintl	127–132

^aRepresentative publications are chosen for each class of materials.

Nevertheless, there is insight in comparing the correlation (i.e., slope) at different temperatures. For example, while there is considerable scatter in the data at each temperature, the temperature dependence of the correlation is strongly linear ($R^2 = 0.997$), in good agreement with the expected behavior as Umklapp scattering begins to dominate at higher temperatures. The changing slope as a function of data set temperature results from increased phonon scattering lowering the phonon

contribution to thermal conductivity at higher temperatures. When the lattice κ_l is minimized, the electronic κ_e dominates, so changes in electrical resistivity have a larger impact on thermal conductivity. The power dependence of temperature on κ vs ρ slope may result from the Umklapp scattering, a three-phonon scattering process. According to the Bose-Einstein expression, $\langle n \rangle = 1/(\exp[\hbar\omega/k_B T] - 1) \approx k_B T/\hbar\omega$, the average phonon population, $\langle n \rangle$, increases with temperature which increases the likelihood of phonon-phonon scattering. In fact, Grimvall⁴⁷ has predicted a T^{-1} temperature dependence for three-phonon scattering at moderate temperatures (near the Debye temperature).

RESOURCE CONSIDERATIONS

In this section we review parameters beyond thermoelectric performance. We analyze resource considerations such as scarcity and supply risk of thermoelectric materials based on their elemental composition. The implications of the analysis on the choice of thermoelectric materials is discussed below. A combined analysis of performance and resource considerations is necessary for identifying thermoelectric materials with the greatest promise of widespread application. Scarce elements could be employed in efforts to make high-impact discoveries and to obtain better understanding of materials trends. However, when materials are meant for widespread deployment, the incorporation of scarce elements becomes an important point to consider during the design phase.

We begin by discussing scarcity which places a fundamental limit on the amount of material available for use. The following section addresses HHI, a measure of geopolitical influence. Abundance of raw materials is a key factor that determines cost, in addition to other factors such as actuarial costs, packaging, transportation, and assembly, to name a few. Technologies relying on scarce elements such as Re, Te, or Pt group metals, are susceptible to large materials costs. For example, turbine engines consume 70% of the world's Re production, which drives the cost of Re.⁴⁸ The cost of scarce materials can be exacerbated when multiple technologies use the same element, such as the use of In for transparent conducting oxides (TCO) and copper indium gallium selenide (CIGS) photovoltaics. Cost can also be high for relatively abundant materials where both the demand and the production are low, such as low purity Si.¹⁹

Plotting scarcity against HHI_p for thermoelectric materials is a useful tool to identify technologies where materials cost can be a concern. About half of the highest efficiency materials contain a large weight fraction of Te (e.g., Be_2Te_3 , PbTe , La_3Te_4 , TAGS, Ag_3TlTe_5 , SbTe , or Tl_3BiTe_6). Given that Te has one of the highest scarcity values, these materials may pose problems for widespread deployment. Te may be rare, but its HHI_p is not critical because it is produced in many countries as a byproduct of Cu, reducing supply risks. Additionally, with proper recycling and waste processing policies, Te could be recovered and reprocessed, a practice already carried out by some photovoltaic companies.

An unanticipated finding from analysis of Figure 8 was that incorporation of even small amounts of scarce elements in otherwise Earth-abundant compounds may be problematic. Mg_2Si has a low scarcity of $\zeta = 27.3$, but when doped with trace amounts of Bi to $\text{Mg}_2\text{Si}_{0.993}\text{Bi}_{0.007}$,³⁷ ζ increases to 1.1×10^6 because of the extremely low crustal abundance (0.025 ppm) and high atomic weight of Bi. Even 0.1% substitution of Bi on the Si site results in a 4 orders of magnitude increase in ζ .

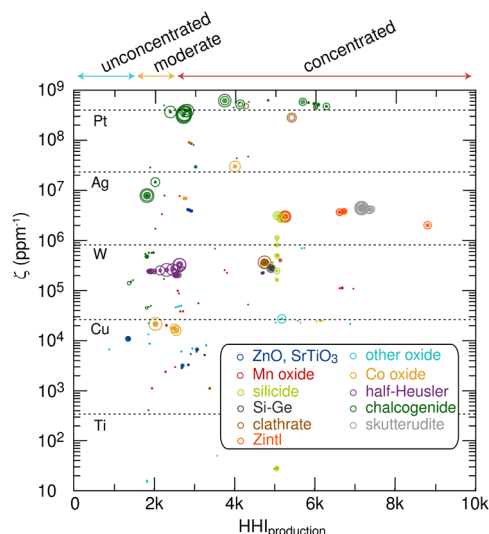


Figure 8. Nearly all high-performance thermoelectric materials pose either a scarcity or HHI risk, which will limit any intensive deployment for energy applications. The scarcity (inverse crustal abundance) of familiar elements is shown to give context. Although not immediately apparent, incorporation of Bi, Sb, Te, or rare-earth elements leads to significant risk in either scarcity or HHI. Materials that pose a lower risk of criticality are SiGe, silicides, half-Heuslers, and cobaltates.

A compromise between performance and resource considerations is seen in oxide materials. Cobalt oxides, such as NaCo_2O_4 , $\text{Ca}_2\text{Co}_2\text{O}_5$, and $\text{Ca}_3\text{Co}_4\text{O}_9$, as well as the reduced ferroelectric oxide $\text{Sr}_{0.61}\text{Ba}_{0.39}\text{Nb}_2\text{O}_6$, have competitive zT values with scarcity values 4 orders of magnitude lower than state-of-the-art chalcogenides, 2 orders of magnitude lower than skutterudites ($\text{ACo}_4\text{Sb}_{12}$, $A = \text{caged atom}$) and Zintl (Zn_4Sb_3 , $\text{Yb}_{14}\text{MnSb}_{11}$), and 1 order of magnitude lower than clathrates ($\text{Ba}_8\text{Ga}_{18}\text{Ge}_{28}$) and half-Heuslers (TiNiSn). ZnO and SrTiO_3 are even more abundant, but have lower performance. It is worth noting that, as with Mg_2Si , if Bi is added to improve performance (e.g., $\text{Bi}_2\text{Sr}_2\text{Co}_2\text{O}_8$), a dramatic increase in ζ is observed.

Determining the Herfindahl-Hirschman Index (HHI) for all elements has shown the following elements to be most at risk for market abuse: Be, K, Br, Rb, Nb, Sb, Cs, W, the rare-earth (RE) elements including Y, and the platinum group elements (PGE): Ru, Rh, Pd, Os, Ir, Pt. These elements have HHI values of greater than 6000 in either production or reserves. As with scarcity, introduction of even small amounts of elements with high HHI_p or HHI_r such as REs, Sb, or Tl, may not be justified for a marginal gain in performance. Materials like Zintl or clathrate antimonides and rare-earth containing chalcogenides are poor outliers compared to other thermoelectric materials in terms of HHI_p risk.

The importance of HHI when considering thermoelectric materials can be illustrated through recent price spikes in antimony. In 2011, a single country that was responsible for nearly 90% of Sb mining halted much of its production. As a result, Sb prices rose over 15% in less than a year. While Sb is particularly concentrated, over 70% of elements exceed the 2500 threshold to be considered "highly concentrated" according to Federal Trade Commission definition. This makes many thermoelectric materials sensitive to changing market conditions. This analysis signals that emphasis should be placed on developing materials with a lower HHI. Examples of materials that have reasonable performance and low HHI

include S- and Se-containing chalcogenides, the half-Heusler materials low in Hf- and Nb-content, and sodium cobalt oxides. Practical consideration of these criticality parameters, although not performance related, would benefit both material and experimental design.

FUTURE DIRECTIONS

Thermoelectric data mining provides insights, enabling rational material design. There are, in addition to the properties described, a number of additional materials properties and parameters that were not included that are nevertheless highly desired in comparing different thermoelectric materials. For example, crystallographic data would allow for the calculation of average atomic volume, unit cell volume, and number of atoms in the unit cell. This could be valuable in looking at changes in thermal conductivity. Mechanical properties, coefficients of thermal expansion, and high temperature stability indices would be valuable in thermoelectric device design. Furthermore, additional parameters not related to performance such as toxicity are of great interest in environmental and life-cycle analyses of thermoelectric components. Some of these properties could be extracted from other databases, measured, or approximated. Other properties, such as toxicity, are fundamentally more challenging to obtain. A comprehensive approach to comparing toxicity across compounds is not well established, since it depends on environmental transformations and persistence, exposure route, form (e.g., oxidation state, presence of counterions), and many other factors. Finally, adding a dimension of time to the plots we have shown may increase their utility. For example, looking at the evolution of HHI over time may provide a better ability to estimate risk and market volatility of particular materials.

The data-driven approach and visualization can be readily extended to other materials functions. For example, photovoltaics, thermal barrier coatings, dielectrics, fuel cells, transparent conducting oxides, gas turbine superalloys, and batteries are all areas of research where progress depends on optimizing several competing requirements concurrently, that could benefit from a comprehensive, data-driven approach.

CLOSING REMARKS

The framework demonstrated here allows researchers to obtain a birds-eye view of a large domain of thermoelectrics research, and compare new materials of interest. The thermoelectric design space is large, and this overview of the field, including important parameters not related to performance, allows researchers to focus on property regions and material families best suited for a given application. In the case of power generation with thermoelectrics, analysis of HHI, scarcity, and materials properties can be done visually. Using the methods shown here, we highlight several families of materials that combine good performance-related properties with noncritical resource availability: highly conductive (electrically) early transition metal oxides, silicides with low bismuth content, and half-Heuslers. We also observe that all high- zT materials are found in the metallic conduction region $0.001 \Omega \text{ cm}$ to $0.01 \Omega \text{ cm}$. This visualization framework is also valuable from a design perspective, where it serves as a useful guide to identifying promising new materials.

AUTHOR INFORMATION

Corresponding Author

*E-mail: mgaultois@mrl.ucsb.edu (M.W.G.), sparks@mrl.ucsb.edu (T.D.S.), seshadri@mrl.ucsb.edu (R.S.).

Notes

The authors declare no competing financial interest.

ACKNOWLEDGMENTS

We thank the National Science Foundation for support of this research through NSF-DMR 1121053. R.S. acknowledges useful discussions with Ambuj Singh. M.W.G. thanks the Natural Sciences and Engineering Council of Canada for support through a NSERC Postgraduate Scholarship and the U.S. Department of State for an International Fulbright Science & Technology Award. The research at Harvard was supported by the Office of Naval Research through contract N000141110894 and by the Harvard University Center for the Environment.

REFERENCES

- (1) Sales, B. C.; Mandrus, D.; Williams, R. K. *Science* **1996**, *272*, 1325–1328.
- (2) Hicks, L. D.; Dresselhaus, M. S. *Phys. Rev. B* **1993**, *47*, 12727–12731.
- (3) Dresselhaus, M. S.; Chen, G.; Tang, M. Y.; Yang, R. G.; Lee, H.; Wang, D. Z.; Ren, Z. F.; Fleurial, J. P.; Gogna, P. *Adv. Mater.* **2007**, *19*, 1043–1053.
- (4) Kaulzarich, S. M.; Brown, S. R.; Snyder, G. J. *Dalton Trans.* **2007**, 2099–2107.
- (5) Pei, Y.; Shi, X.; LaLonde, A.; Wang, H.; Chen, L.; Snyder, G. J. *Nature* **2011**, *473*, 66–69.
- (6) Hsu, K. F.; Loo, S.; Guo, F.; Chen, W.; Dyck, J. S.; Uher, C.; Hogan, T.; Polychroniadis, E. K.; Kanatzidis, M. G. *Science* **2004**, *303*, 818–821.
- (7) Biswas, K.; Good, M. S.; Roberts, K. C.; Subramanian, M. A.; Hendricks, T. J. *J. Mater. Res.* **2011**, *26*, 1827–1835.
- (8) Biswas, K.; He, J.; Blum, I. D.; Wu, C.-I.; Hogan, T. P.; Seidman, D. N.; Draid, V. P.; Kanatzidis, M. G. *Nature* **2012**, *489*, 414–418.
- (9) Snyder, G. J.; Toberer, E. S. *Nat. Mater.* **2008**, *7*, 105–114.
- (10) Nolas, G. S.; Sharp, J.; Goldsmid, H. J. *Thermoelectrics: Basic principles and new materials developments*; Springer series in materials science 45; Springer: Berlin, Germany, 2001.
- (11) Goldsmid, H. J. *Introduction to thermoelectricity*; Springer: Heidelberg, Germany, 2009.
- (12) Rowe, D. M. *CRC Handbook of thermoelectrics*; CRC Press: Boca Raton, FL, 1995.
- (13) Rowe, D. M. *Thermoelectrics handbook: Macro to nano*; CRC Press: Boca Raton, FL, 2006.
- (14) Bucholz, E. W.; Kong, C. S.; Marchman, K. R.; Sawyer, W. G.; Phillpot, S. R.; Sinnott, S. B.; Rajan, K. *Tribol. Lett.* **2012**, *47*, 211–221.
- (15) Kong, C. S.; Villars, P.; Iwata, S.; Rajan, K. *Comput. Sci. Discovery* **2012**, *5*, 015004.
- (16) Curtarolo, S.; Hart, G. L. W.; Nardelli, M. B.; Mingo, N.; Sanvito, S.; Levy, O. *Nat. Mater.* **2013**, *12*, 191–201.
- (17) Jonker, G. H. *Philips Res. Rep.* **1968**, *23*, 131–138.
- (18) Graedel, T. E. *Annu. Rev. Mater. Res.* **2011**, *41*, 323–335.
- (19) Himm, G.; Klar, P. J. *Status Solidi RRL* **2011**, *5*, 324–331.
- (20) Herfindahl, O. C. Concentration in the US Steel Industry. Ph.D. Thesis, Columbia University, New York, 1950.
- (21) Hirschman, A. O. *National power and the structure of foreign trade*; California University Press: Berkeley, CA, 1980.
- (22) Rosenau-Tornow, D.; Buchholz, P.; Riemann, A.; Wagner, M. *Resour. Policy* **2009**, *34*, 161–175.
- (23) *Fact Sheet: Rare Earth Oxides (REO)*; POLINARES Working Paper, 2012; Vol. 37.

- (24) Huwaldt, J. A. *Plot Digitizer*; 2011; <http://plotdigitizer.sourceforge.net/>.
- (25) Tummers, B. *DataThief III*; 2006; <http://datathief.org/>.
- (26) U.S. Department of Justice and the Federal Trade Commission. *Horizontal merger guidelines*; 2010.
- (27) Haynes, W., Ed.; *CRC Handbook of chemistry and physics*, 93rd ed.; CRC Press: Boca Raton, FL, 2012.
- (28) Gaultois, M. W.; Barton, P. T.; Birkel, C. S.; Misch, L. M.; Rodriguez, E. E.; Stucky, G. D.; Seshadri, R. *J. Phys.: Condens. Matter* **2013**, *25*, 186004.
- (29) Bak, T.; Nowotny, J. *J. Phys. Chem. C* **2011**, *115*, 9746–9752.
- (30) Heremans, J. P.; Jovovic, V.; Toberer, E. S.; Saramat, A.; Kurosaki, K.; Charoenphakdee, A.; Yamanaka, S.; Snyder, G. J. *Science* **2008**, *321*, 554–557.
- (31) Koshibae, W.; Tsutsui, K.; Maekawa, S. *Phys. Rev. B* **2000**, *62*, 6869–6872.
- (32) Singh, D. J.; Kasinathan, D. *J. Electron. Mater.* **2007**, *36*, 736–739.
- (33) Tsubota, T.; Ohtaki, M.; Eguchi, K.; Arai, H. *J. Mater. Chem.* **1997**, *7*, 85–90.
- (34) Fujita, K.; Mochida, T.; Nakamura, K. *Jpn. J. Appl. Phys., Part 1* **2001**, *40*, 4644–4647.
- (35) Funahashi, R.; Matsubara, I.; Ikuta, H.; Takeuchi, T.; Mizutani, U.; Sodeoka, S. *Jpn. J. Appl. Phys., Part 2* **2000**, *39*, L1127–L1129.
- (36) Sakurada, S.; Shutoh, N. *Appl. Phys. Lett.* **2005**, *86*, 082105.
- (37) Bux, S. K.; Yeung, M. T.; Toberer, E. S.; Snyder, G. J.; Kaner, R. B.; Fleuriel, J. P. *J. Mater. Chem.* **2011**, *21*, 12259–12266.
- (38) Fukano, M.; Iida, T.; Makino, K.; Akasaka, M.; Oguni, Y.; Takanashi, Y. *MRS Proc.* **2007**, *1044*, 1044–U06–13.
- (39) Vining, C. B.; Laskow, W.; Hanson, J. O.; Vanderbeck, R. R.; Gorsuch, P. D. *J. Appl. Phys.* **1991**, *69*, 4333–4340.
- (40) Chung, D. Y.; Hogan, T.; Brazis, P.; Rocci-Lane, M.; Kannewurf, C.; Bastea, M.; Uher, C.; Kanatzidis, M. G. *Science* **2000**, *287*, 1024–1027.
- (41) Anno, H.; Hokazono, M.; Kawamura, M.; Nagao, J.; Matsubara, K. In *Proceedings of the 21st International Conference on Thermoelectrics*, Long Beach, CA, August 25–29, 2002; IEEE: New York, 2002; pp 77–80.
- (42) Toberer, E. S.; Christensen, M.; Iversen, B. B.; Snyder, G. J. *Phys. Rev. B* **2008**, *77*, 075203.
- (43) He, T.; Chen, J. Z.; Rosenfeld, H. D.; Subramanian, M. A. *Chem. Mater.* **2006**, *18*, 759–762.
- (44) Funahashi, R.; Mikami, M.; Urata, S.; Kitawaki, M.; Kouuchi, T.; Mizuno, K. *Meas. Sci. Technol.* **2005**, *16*, 70–80.
- (45) DiSalvo, F. J. *Science* **1999**, *285*, 703–706.
- (46) Clarke, D. R. *Surf. Coat. Technol.* **2003**, *163*, 67–74.
- (47) Grimvall, G. *Thermophysical properties of materials*, 1st ed.; Elsevier: Amsterdam, The Netherlands, 1999.
- (48) Naumov, A. *Russ. J. Non-Ferrous Met.* **2007**, *48*, 418–423.
- (49) Lan, J.; Lin, Y.; Mei, A.; Nan, C.; Liu, Y.; Zhang, B.; Li, J. *J. Mater. Sci. Technol.* **2009**, *25*, 535–538.
- (50) Park, J. W.; Kwak, D. H.; Yoon, S. H.; Choi, S. C. *J. Alloys Compd.* **2009**, *487*, 550–555.
- (51) Populoh, S.; Trottmann, M.; Aguire, M. H.; Weidenkaff, A. *J. Mater. Res.* **2011**, *26*, 1947–1952.
- (52) Zhou, Y. Q.; Matsubara, I.; Funahashi, R.; Xu, G. J.; Shikano, M. *Mater. Res. Bull.* **2003**, *38*, 341–346.
- (53) Lan, J. L.; Lin, Y. H.; Fang, H.; Mei, A.; Nan, C. W.; Liu, Y.; Xu, S. L.; Peters, M. *J. Am. Ceram. Soc.* **2010**, *93*, 2121–2124.
- (54) Flahaut, D.; Mihara, T.; Funahashi, R.; Nabeshima, N.; Lee, K.; Ohta, H.; Koumoto, K. *J. Appl. Phys.* **2006**, *100*, 084911.
- (55) Kobayashi, T.; Takizawa, H.; Endo, T.; Sato, T.; Shimada, M.; Taguchi, H.; Nagao, M. *J. Solid State Chem.* **1991**, *92*, 116–129.
- (56) Sparks, T. D.; Gurlo, A.; Clarke, D. R. *J. Mater. Chem.* **2012**, *22*, 4631–4636.
- (57) Flahaut, D.; Funahashi, R.; Lee, K.; Ohta, H.; Koumoto, K. In *Proceedings of the 25th International Conference on Thermoelectrics*, Vienna, Austria, August 6–10, 2006; IEEE: New York, 2006; pp 103–106.
- (58) Hejtmanek, J.; Jirak, Z.; Marysko, M.; Martin, C.; Maignan, A.; Hervieu, M.; Raveau, B. *Phys. Rev. B* **1999**, *60*, 14057–14065.
- (59) Kobayashi, W.; Terasaki, I.; Mikami, M.; Funahashi, R.; Nomura, T.; Katsufuji, T. *J. Appl. Phys.* **2004**, *95*, 6825–6827.
- (60) Ohtaki, M.; Koga, H.; Tokunaga, T.; Eguchi, K.; Arai, H. *J. Solid State Chem.* **1995**, *120*, 105–111.
- (61) Kim, K. H.; Shim, S. H.; Shim, K. B.; Niihara, K.; Hojo, J. *J. Am. Ceram. Soc.* **2005**, *88*, 628–632.
- (62) Berardan, D.; Guilmeau, E.; Maignan, A.; Raveau, B. *Solid State Commun.* **2008**, *146*, 97–101.
- (63) Guilmeau, E.; Maignan, A.; Martin, C. *J. Electron. Mater.* **2009**, *38*, 1104–1108.
- (64) Hebert, S.; Flahaut, D.; Martin, C.; Lemonnier, S.; Noudem, J.; Goupil, C.; Maignan, A.; Hejtmanek, J. *Prog. Solid State Chem.* **2007**, *35*, 457–467.
- (65) Wang, Y.; Lee, K. H.; Ohta, H.; Koumoto, K. *J. Appl. Phys.* **2009**, *105*, 103701.
- (66) Muta, H.; Kurosaki, K.; Yamanaka, S. *J. Alloys Compd.* **2003**, *350*, 292–295.
- (67) Park, K.; Ko, K. Y.; Seo, W. S.; Cho, W. S.; Kim, J. G.; Kim, J. Y. *J. Eur. Ceram. Soc.* **2007**, *27*, 813–817.
- (68) Sonne, M.; van Nong, N.; He, Z.; Pryds, N.; Linderoth, S. In *Proceedings of the 8th European Conference on Thermoelectrics*, Como, Italy, September 22–24, 2010; International Thermoelectric Society: Pasadena, CA, 2010.
- (69) Wang, N.; He, H. C.; Li, X. A.; Han, L.; Zhang, C. Q. *J. Alloys Compd.* **2010**, *506*, 293–296.
- (70) Obara, H.; Yamamoto, A.; Lee, C. H.; Kobayashi, K.; Matsumoto, A.; Funahashi, R. *Jpn. J. Appl. Phys., Part 2* **2004**, *43*, L540–L542.
- (71) Muta, H.; Kurosaki, K.; Yamanaka, S. *J. Alloys Compd.* **2004**, *368*, 22–24.
- (72) Jood, P.; Mehta, R. J.; Zhang, Y.; Peleckis, G.; Wang, X.; Siegel, R. W.; Borca-Tasciuc, T.; Dou, S. X.; Ramanath, G. *Nano Lett.* **2011**, *11*, 4337–4342.
- (73) Funahashi, R.; Shikano, M. *Appl. Phys. Lett.* **2002**, *81*, 1459–1461.
- (74) Iwasaki, K.; Ito, T.; Nagasaki, T.; Arita, Y.; Yoshino, M.; Matsui, T. *J. Solid State Chem.* **2008**, *181*, 3145–3150.
- (75) Nan, J.; Wu, J.; Deng, Y.; Nan, C. W. *Solid State Commun.* **2002**, *124*, 243–246.
- (76) Shikano, M.; Funahashi, R. *Appl. Phys. Lett.* **2003**, *82*, 1851–1853.
- (77) Xu, G. J.; Funahashi, R.; Shikano, M.; Matsubara, I.; Zhou, Y. Q. *Appl. Phys. Lett.* **2002**, *80*, 3760–3762.
- (78) Shin, W.; Murayama, N. *Jpn. J. Appl. Phys., Part 2* **1999**, *38*, L1336–L1338.
- (79) Zhou, Y. Q.; Matsubara, I.; Funahashi, R.; Sodeoka, S. *Mater. Lett.* **2001**, *51*, 347–350.
- (80) Yasukawa, M.; Kuniyoshi, S.; Kono, T. *Solid State Commun.* **2003**, *126*, 213–216.
- (81) Yasukawa, M.; Murayama, N. *J. Mater. Sci. Lett.* **1997**, *16*, 1731–1734.
- (82) Funahashi, R.; Mikami, M.; Urata, S.; Kouuchi, T.; Mizuno, K.; Chong, K. *MRS Proc.* **2003**, *793*, S3.3.
- (83) Sanmathi, C. S.; Retoux, R.; Noudem, J. G. In *Proceedings of the 26th International Conference on Thermoelectrics*, Jeju, Korea, June 3–7, 2007; IEEE: New York, 2007; pp 171–174.
- (84) Choi, M. Y.; Kim, J. S. *Phys. Rev. B* **1999**, *59*, 192–194.
- (85) Wang, C. X.; Chen, X. B.; Zhu, A. P. *Solid State Commun.* **2012**, *152*, 1067–1071.
- (86) Takagi, H.; Batlogg, B.; Kao, H. L.; Kwo, J.; Cava, R. J.; Krajewski, J. J.; Peck, W. F. *Phys. Rev. Lett.* **1992**, *69*, 2975–2978.
- (87) Dong, X.; Wang, H.; Hua, Z.; Peng, S.; Dong, L.; Wang, Y. *J. Mater. Sci.: Mater. Electron.* **2012**, *23*, 1210–1214.
- (88) Kuriyama, H.; Nohara, M.; Sasagawa, T.; Takubo, K.; Mizokawa, T.; Kimura, K.; Takagi, H. In *Proceedings of the 25th International Conference on Thermoelectrics*, Vienna, Austria, August 6–10, 2006; IEEE: New York, 2006; pp 97–98.

- (89) Ono, Y.; Satoh, K.; Nozaki, T.; Kajitani, T. In *Proceedings of the 25th International Conference on Thermoelectrics*, Vienna, Austria, August 6–10, 2006; IEEE: New York, 2006; pp 92–96.
- (90) Tinh, N.; Tsuji, T. In *Physics and Engineering of New Materials*; Cat, D., Pucci, A., Wandelt, K., Eds.; Springer: Berlin, Germany, 2009; Vol. 127; pp 209–217.
- (91) Maignan, A.; Eyert, V.; Martin, C.; Kremer, S.; Fresard, R.; Pelloquin, D. *Phys. Rev. B* **2009**, *80*, 115103.
- (92) Muta, H.; Kurosaki, K.; Uno, M.; Yamanaka, S. *J. Alloys Compd.* **2002**, *335*, 200–202.
- (93) Kurosaki, K.; Kobayashi, H.; Yamanaka, S. *J. Alloys Compd.* **2003**, *350*, 340–343.
- (94) Lee, S.; Wilke, R. H. T.; Trolrier-McKinstry, S.; Zhang, S.; Randall, C. A. *Appl. Phys. Lett.* **2010**, *96*, 031910.
- (95) Lee, S.; Dursun, S.; Duran, C.; Randall, C. A. *J. Mater. Res.* **2011**, *26*, 26–30.
- (96) Sakai, A.; Kanno, T.; Takahashi, K.; Omote, A.; Adachi, H.; Yamada, Y. *J. Am. Ceram. Soc.* **2012**, *95*, 1750–1755.
- (97) Joshi, G.; Lee, H.; Lan, Y.; Wang, X.; Zhu, G.; Wang, D.; Gould, R. W.; Cuff, D. C.; Tang, M. Y.; Dresselhaus, M. S.; Chen, G.; Ren, Z. *Nano Lett.* **2008**, *8*, 4670–4674.
- (98) Saramat, A.; Svensson, G.; Palmqvist, A. E. C.; Stiewe, C.; Mueller, E.; Platzek, D.; Williams, S. G. K.; Rowe, D. M.; Bryan, J. D.; Stucky, G. D. *J. Appl. Phys.* **2006**, *99*, 023708.
- (99) Wolfing, B.; Kloc, C.; Teubner, J.; Bucher, E. *Phys. Rev. Lett.* **2001**, *86*, 4350–4353.
- (100) Candolfi, C.; Aydemir, U.; Baitinger, M.; Oeschler, N.; Steglich, F.; Grin, Y. *J. Appl. Phys.* **2012**, *111*, 043706.
- (101) Nolas, G. S.; Cohn, J. L.; Slack, G. A.; Schujman, S. B. *Appl. Phys. Lett.* **1998**, *73*, 178–180.
- (102) Kuznetsov, V. I.; Kuznetsova, L. A.; Kaliazin, A. E.; Rowe, D. M. In *Proceedings of the 18th International Conference on Thermoelectrics*, Baltimore, MD, Aug 29–Sept 2, 1999; IEEE: New York, 1999; pp 177–180.
- (103) Roudebush, J. H.; Toberer, E. S.; Hope, H.; Snyder, G. J.; Kauzlarich, S. M. *J. Solid State Chem.* **2011**, *184*, 1176–1185.
- (104) Katsuyama, S.; Matsushima, H.; Ito, M. *J. Alloys Compd.* **2004**, *385*, 232–237.
- (105) Kimura, Y.; Tamura, Y.; Kita, T. *Appl. Phys. Lett.* **2008**, *92*, 012105.
- (106) Muta, H.; Kanemitsu, T.; Kurosaki, K.; Yamanaka, S. *J. Alloys Compd.* **2009**, *469*, 50–55.
- (107) Shutoh, N.; Sakurada, S. *J. Alloys Compd.* **2005**, *389*, 204–208.
- (108) Fleurial, J. P.; Borshchevsky, A.; Caillat, T.; Morelli, D. T.; Meisner, G. P. In *Proceedings of the 15th International Conference on Thermoelectrics* Pasadena, CA, March 25–28, 1996; IEEE: New York, 1996; pp 91–95.
- (109) Scherrer, H.; Scherrer, S. In *CRC Handbook of thermoelectrics*; Rowe, D. M., Ed.; CRC Press: Boca Raton, FL, 1995.
- (110) Skrabek, E. A.; Trimmer, D. S. In *CRC Handbook of thermoelectrics*; Rowe, D. M., Ed.; CRC Press: Boca Raton, FL, 1995.
- (111) Chung, D. Y.; Choi, K. S.; Iordanidis, L.; Schindler, J. L.; Brazis, P. W.; Kannewurf, C. R.; Chen, B. X.; Hu, S. Q.; Uher, C.; Kanatzidis, M. G. *Chem. Mater.* **1997**, *9*, 3060–3071.
- (112) Gascoin, F.; Maignan, A. *Chem. Mater.* **2011**, *23*, 2510–2513.
- (113) Kurosaki, K.; Kosuga, A.; Yamanaka, S. *J. Alloys Compd.* **2003**, *351*, 208–211.
- (114) Ohta, M.; Yamamoto, A.; Obara, H. *J. Electron. Mater.* **2010**, *39*, 2117–2121.
- (115) Wan, C. L.; Wang, Y. F.; Wang, N.; Koumoto, K. *Materials* **2010**, *3*, 2606–2617.
- (116) Kanatzidis, M. G.; McCarthy, T. J.; Tanzer, T. A.; Chen, L. H.; Iordanidis, L.; Hogan, T.; Kannewurf, C. R.; Uher, C.; Chen, B. X. *Chem. Mater.* **1996**, *8*, 1465–1474.
- (117) McGuire, M. A.; Reynolds, T. K.; DiSalvo, F. J. *Chem. Mater.* **2005**, *17*, 2875–2884.
- (118) Kurosaki, K.; Kosuga, A.; Muta, H.; Uno, M.; Yamanaka, S. *Appl. Phys. Lett.* **2005**, *87*, 061919.
- (119) Sharp, J. W.; Sales, B. C.; Mandrus, D. G.; Chakoumakos, B. C. *Appl. Phys. Lett.* **1999**, *74*, 3794–3796.
- (120) Liu, H. L.; Shi, X.; Xu, F. F.; Zhang, L. L.; Zhang, W. Q.; Chen, L. D.; Li, Q.; Uher, C.; Day, T.; Snyder, G. J. *Nat. Mater.* **2012**, *11*, 422–425.
- (121) May, A. F.; Flage-Larsen, E.; Snyder, G. J. *Phys. Rev. B* **2010**, *81*, 125205.
- (122) Raghavendra, N.; Gascoin, F.; Takahashi, H.; Hebert, S.; Guilmeau, E.; Terasaki, I. Thermoelectric Properties of Hollandite Chalcogenides AM_3X_8 . Presented at EMRS 2012 Spring Meeting, Strasbourg, France, May 15–17.
- (123) Mars, K.; Ihou-Mouko, H.; Pont, G.; Tobola, J.; Scherrer, H. *J. Electron. Mater.* **2009**, *38*, 1360–1364.
- (124) Zhang, Q.; Zhao, X. B.; Yin, H.; Zhu, T. *J. Alloys Compd.* **2008**, *464*, 9–12.
- (125) Ito, M.; Tada, T.; Hara, S. *J. Alloys Compd.* **2006**, *408–412*, 363–367.
- (126) Akasaka, M.; Iida, T.; Nemoto, T.; Soga, J.; Sato, J.; Makino, K.; Fukano, M.; Takanashi, Y. *J. Cryst. Growth* **2007**, *304*, 196–201.
- (127) Brown, S. R.; Kauzlarich, S. M.; Gascoin, F.; Snyder, G. J. *Chem. Mater.* **2006**, *18*, 1873–1877.
- (128) Caillat, T.; Fleurial, J. P.; Borshchevsky, A. *J. Phys. Chem. Solids* **1997**, *58*, 1119–1125.
- (129) Kitagawa, H.; Noguchi, H.; Kiyabu, T.; Itoh, M.; Noda, Y. *J. Phys. Chem. Solids* **2004**, *65*, 1223–1227.
- (130) Toberer, E. S.; Rauwel, P.; Gariel, S.; Taftø, J.; Snyder, G. J. *J. Mater. Chem.* **2010**, *20*, 9877–9885.
- (131) Zevalkink, A.; Toberer, E. S.; Zeier, W. G.; Flage-Larsen, E.; Snyder, G. J. *Energy Environ. Sci.* **2011**, *4*, 510–518.
- (132) Toberer, E. S.; Zevalkink, A.; Crisosto, N.; Snyder, G. J. *Adv. Funct. Mater.* **2010**, *20*, 4375–4380.